

Chiral extrapolations of lattice QCD results

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Outline

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Masses and decay constants

Form factors

Scattering amplitudes

Summary and conclusions

Chiral extrapolations: an opportunity

- ▶ Extrapolations (the continuum, the chiral and the infinite volume) are rather an opportunity than a nuisance
- ▶ Especially the chiral extrapolation allows us to learn something about QCD which we can not get from phenomenology
- ▶ In doing this we also check the lattice calculations against some firm predictions of CHPT (the chiral logs)

Chiral extrapolations: an opportunity

- ▶ Extrapolations (the continuum, the chiral and the infinite volume) are rather an opportunity than a nuisance
- ▶ Especially the chiral extrapolation allows us to learn something about QCD which we can not get from phenomenology
- ▶ In doing this we also check the lattice calculations against some firm predictions of CHPT (the chiral logs)

But

- ▶ the chiral region may be rather small (we start entering it now)
- ▶ the necessary precision and control over systematic effects may be very high

Pion mass and decay constant

$$M_\pi^2 = M^2 \left\{ 1 - \frac{1}{2} x \ln \frac{\Lambda_3^2}{M^2} + \frac{17}{8} x^2 \left(\ln \frac{\Lambda_M^2}{M^2} \right)^2 + x^2 k_M + O(x^3) \right\}$$

$$F_\pi = F \left\{ 1 + x \ln \frac{\Lambda_4^2}{M^2} - \frac{5}{4} x^2 \left(\ln \frac{\Lambda_F^2}{M^2} \right)^2 + x^2 k_F + O(x^3) \right\}$$

$$\text{where } x = \frac{M^2}{(4\pi F)^2} \quad M^2 = 2Bm$$

$\Lambda_{3,4}$ are low energy constants (LEC) in form of energy scales:

$$\bar{\ell}_n = \ln \frac{\Lambda_n^2}{M_\pi^2} \quad n = 1, \dots, 7$$

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where $x = \frac{M^2}{(4\pi F)^2}$ $M^2 = 2Bm$

$\Lambda_{M,F}$ are combinations thereof

$$\ln \frac{\Lambda_M^2}{M^2} = \frac{1}{51} \left(28 \ln \frac{\Lambda_1^2}{M^2} + 32 \ln \frac{\Lambda_2^2}{M^2} - 9 \ln \frac{\Lambda_3^2}{M^2} + 49 \right)$$

$$\ln \frac{\Lambda_F^2}{M^2} = \frac{1}{30} \left(14 \ln \frac{\Lambda_1^2}{M^2} + 16 \ln \frac{\Lambda_2^2}{M^2} + 6 \ln \frac{\Lambda_3^2}{M^2} - 6 \ln \frac{\Lambda_4^2}{M^2} + 23 \right)$$

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where $x = \frac{M^2}{(4\pi F)^2}$ $M^2 = 2Bm$

and $k_{M,F} O(p^6)$ LEC

Pion mass and decay constant – Numerics

Numbers from the phenomenology: $F = 86.2 \pm 0.5 \text{ MeV}$

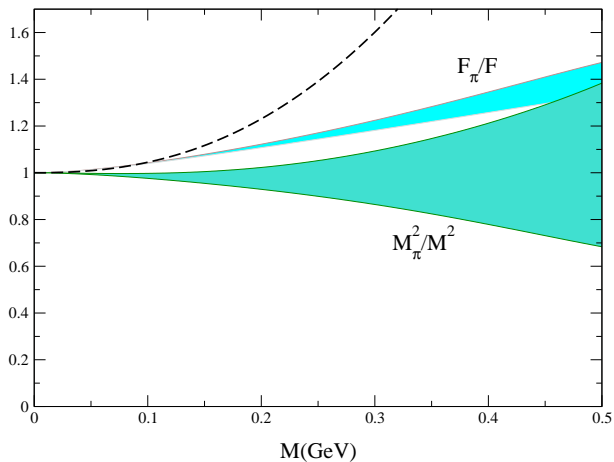
$$\begin{aligned} \bar{l}_1 = -0.4 \pm 0.6 & \Leftrightarrow \Lambda_1 = 0.12^{+0.04}_{-0.03} \text{ GeV} \\ \bar{l}_2 = 4.3 \pm 0.1 & \Leftrightarrow \Lambda_2 = 1.2 \pm 0.06 \text{ GeV} \\ \bar{l}_3 = 2.9 \pm 2.4 & \Leftrightarrow \Lambda_3 = 0.6^{+1.4}_{-0.4} \text{ GeV} \\ \bar{l}_4 = 4.4 \pm 0.2 & \Leftrightarrow \Lambda_4 = 1.25^{+0.15}_{-0.13} \text{ GeV} \end{aligned}$$

From these, it follows:

$$\Lambda_M = 0.6^{+0.18}_{-0.15} \text{ GeV} \quad \Lambda_F = 0.5^{+0.16}_{-0.13} \text{ GeV}$$

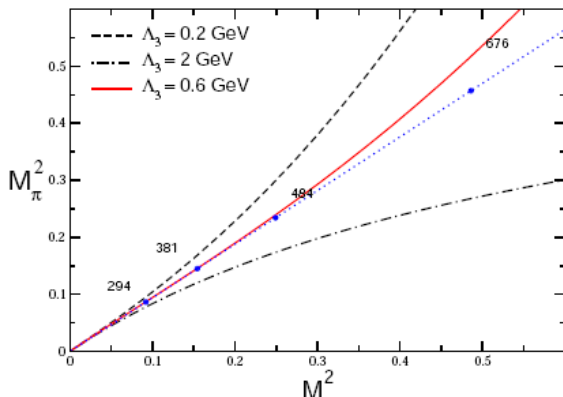
The NNLO LEC's k_F and k_M can only be roughly estimated

Pion mass and decay constant – Plots



From CG, Gasser, Leutwyler (01)

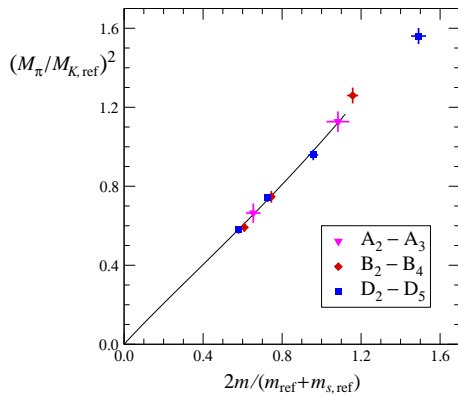
Pion mass and decay constant – Plots



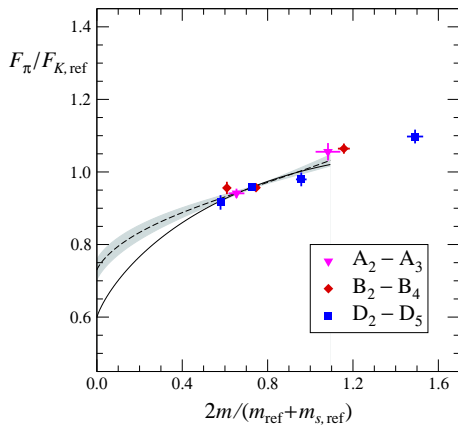
Lattice data from Del Debbio et al.

Thanks to H. Leutwyler for the figure

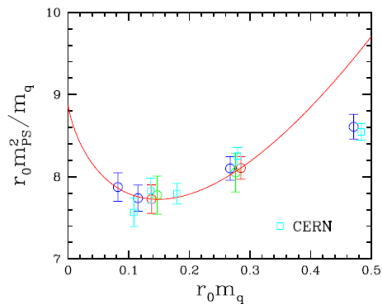
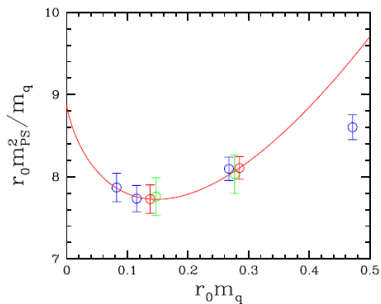
Recent lattice results



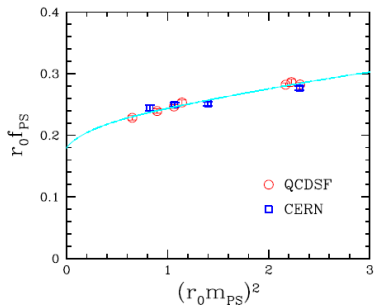
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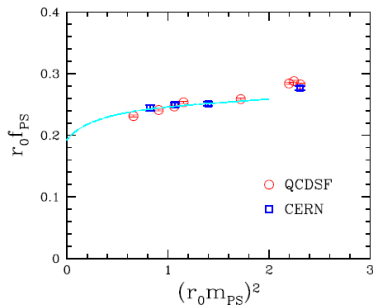
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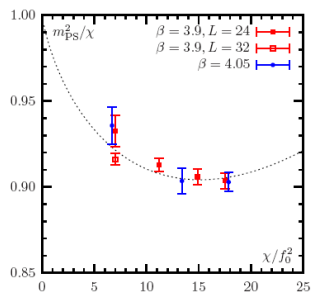
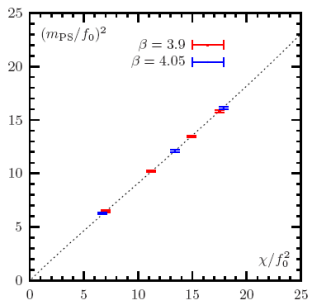


$$r_0 f_0 = 0.179(2) \quad r_0 \Lambda_4 = 3.32(6)$$



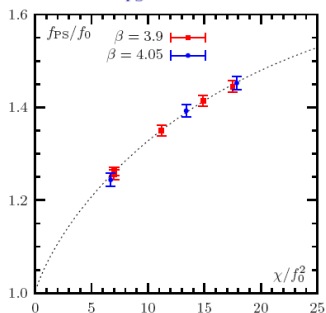
$$r_0 f_0 = 0.192(3) \quad r_0 \Lambda_4 = 3.32(6)$$

Recent lattice results



Recent lattice results

Pion-Sector: f_{PS} as function of the Quark Mass



- fit: $\chi^2/\text{dof} = 1.2$
- lattice spacings:

$$a(3.90) = 0.0855(5)(3) \text{ fm}$$

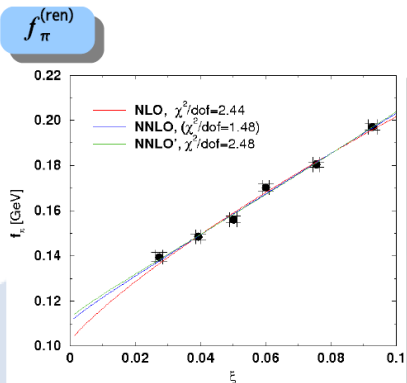
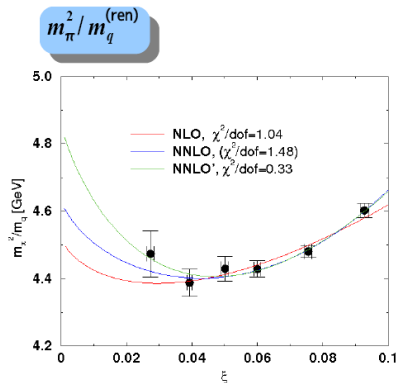
$$a(4.05) = 0.0666(6)(9) \text{ fm}$$

using $f_\pi = 130.7 \text{ MeV}$ and

$$m_\pi^0 = 135 \text{ MeV} \quad [\text{Aubin et al., 2004}]$$

$$\rightarrow r_0 = 0.441(14)(5) \text{ fm}$$

Recent lattice results



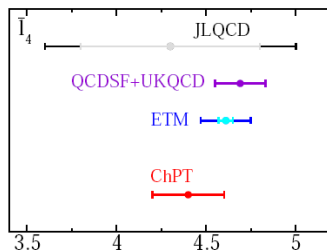
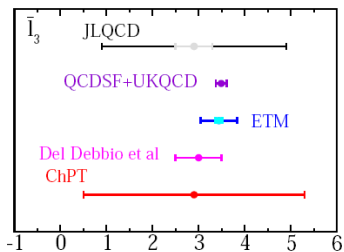
Summary of recent lattice results

Summary of the lattice results from recent $N_f = 2$ simulations ('06,'07):

	F (MeV)	$\Sigma^{\overline{\text{MS}}}(2 \text{ GeV})$ (MeV) ³	$\bar{\Gamma}_3$	$\bar{\Gamma}_4$
Del Debbio et al			3.0(5)	
ETM	85.98(7)(21)(35)	266(6)(0)(6)	3.44(8)(26)(6)	4.61(4)(3)(7)
QCDSF/UKQCD	79(5)	273(12)	3.49(12)	4.69(14)
JLQCD	78.1(2.7)(1.2)	242(6)(6)	2.9(4)(1.6)	4.3(5)(2)

Summary of recent lattice results

Results for \bar{T}_3, \bar{T}_4 from $N_f = 2$ simulations:



Chiral predictions for a_0^0 and a_0^2

$$\begin{aligned}
 a_0^0 &= 0.220 \pm 0.001 + 0.027\Delta_{r^2} - 0.0017\Delta\ell_3 \\
 10 \cdot a_0^2 &= -0.444 \pm 0.003 - 0.04\Delta_{r^2} - 0.004\Delta\ell_3
 \end{aligned}$$

where

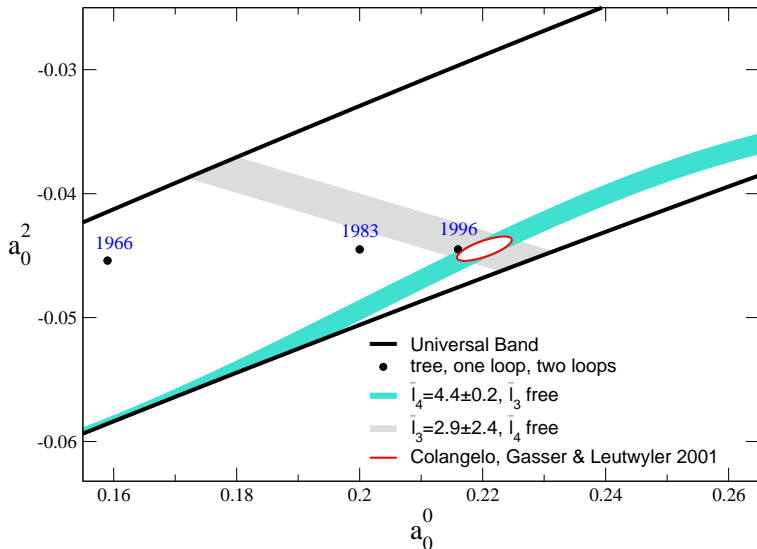
$$\langle r^2 \rangle_s = 0.61 \text{fm}^2 (1 + \Delta_{r^2}) \quad \bar{\ell}_3 = 2.9 + \Delta\ell_3$$

$$\text{with } \langle r^2 \rangle_s = \frac{3}{8\pi^2 F_\pi^2} \left\{ \bar{\ell}_4 - \frac{13}{12} + \xi \Delta_r + \mathcal{O}(\xi^2) \right\} \quad \xi \equiv \left(\frac{M_\pi^2}{4\pi F_\pi^2} \right)$$

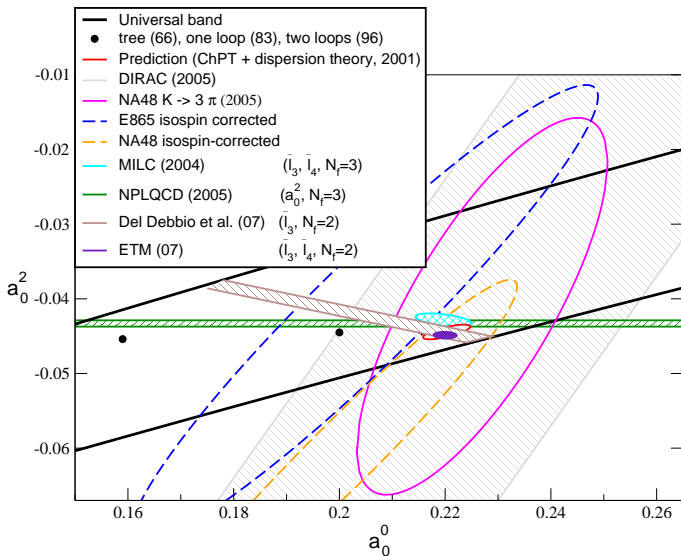
Adding errors in quadrature

$$[\Delta_{r^2} = 6.5\%, \Delta\ell_3 = 2.4]$$

$$\begin{aligned}
 a_0^0 &= 0.220 \pm 0.005 \\
 10 \cdot a_0^2 &= -0.444 \pm 0.01 \\
 a_0^0 - a_0^2 &= 0.265 \pm 0.004
 \end{aligned}$$

Chiral predictions for a_0^0 and a_0^2 

The $\pi\pi$ S-wave scattering lengths plane



Sensitivity to the quark condensate

The constant $\bar{\ell}_3$ determines the NLO quark mass dependence of the pion mass

$$M_\pi^2 = 2B\hat{m} \left[1 + \frac{2B\hat{m}}{16\pi F_\pi^2} \bar{\ell}_3 + \mathcal{O}(\hat{m}^2) \right]$$
$$\hat{m} = \frac{m_u + m_d}{2} \quad B = -\frac{1}{F^2} \langle 0 | \bar{q}q | 0 \rangle$$

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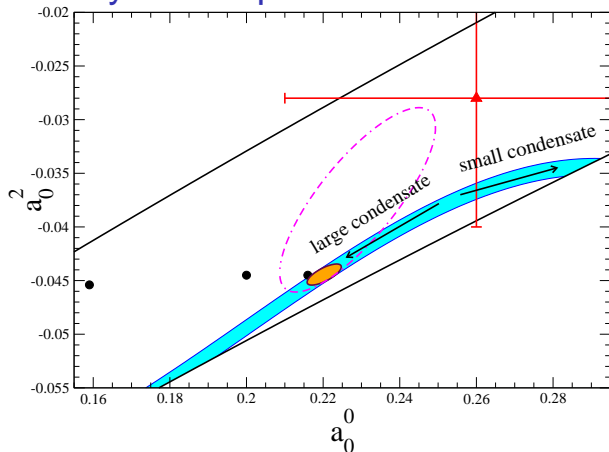
$$\hat{m} = \frac{m_u + m_d}{2} \quad B = -\frac{1}{F^2} \langle 0 | \bar{q}q | 0 \rangle$$

Its size tells us what fraction of the pion mass is given by the Gell-Mann–Oakes–Renner term

$$M_{\text{GMOR}}^2 \equiv 2B\hat{m}$$

or how large is the quark condensate, the order parameter of chiral symmetry breaking – as emphasized since long by **Jan Stern and collaborators**

Sensitivity to the quark condensate



The E865 data on $K_{\ell 4}$ imply that

GC, Gasser and Leutwyler PRL (01)

$$M_{\text{GMOR}} > 94\% M_{\pi}$$

Recently confirmed by NA48 data on K_{e4} and $K \rightarrow 3\pi$

Masses and decay constants in $SU(3)$

$$M_K^2 = M_K^{02} \left\{ 1 + \frac{2}{3} \mu_\eta + \frac{8M_K^{02}}{F_0^2} [(2L_8^r - L_5^r) + (2 + x_{\pi K})(2L_6^r - L_4^r)] + \mathcal{O}(M^4) \right\}$$

$$F_K = F_0 \left\{ 1 - \frac{3}{4} \mu_\pi - \frac{3}{2} \mu_K - \frac{3}{4} \mu_\eta + \frac{4M_K^{02}}{F_0^2} [L_5^r + (2 + x_{\pi K})L_4^r] + \mathcal{O}(M^4) \right\}$$

where

$$x_{\pi K} = \left(\frac{M_\pi^0}{M_K^0} \right)^2 \quad \mu_P = \frac{M_P^2}{32\pi^2 F_P^2} \ln \frac{M_P^2}{\mu^2}$$

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Two-loop formulae for the $SU(3)$ masses and decay constants are also available

Masses and decay constants in $SU(3)$ – Numerics

New preliminary results from MILC, RBC-UKQCD and PACS-CS: L_i at the scale M_ρ

LEC $\cdot 10^3$	MILC	RBC-UKQCD	PACS-CS
$(2L_8 - L_5)$	0.3(1)(1)	0.247(45)	-0.23(5)
$(2L_6 - L_4)$	0.3(1) $\binom{+2}{-3}$	-0.02(42)	0.10(4)
L_4	0.1(3) $\binom{+3}{-1}$	0.136(80)	-0.02(11)
L_5	1.4(2) $\binom{+2}{-1}$	0.862(99)	1.47(15)

Talk by S. Necco at Lattice 07

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cf. fits of Bijmans et al. (2000)

$$[L_4^r \equiv L_6^r \equiv 0]$$

$$\mathcal{O}(p^4) : \quad 2L_8^r - L_5^r = 0.5 \cdot 10^{-3}$$

$$L_8^r = 1.0 \cdot 10^{-3}$$

$$\mathcal{O}(p^6) : \quad L_5^r = (0.97 \pm 0.11) \cdot 10^{-3}$$

$$L_8^r = (0.60 \pm 0.18) \cdot 10^{-3}$$

Masses and decay constants in $SU(3)$ – Numerics

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Talk by S. Necco at Lattice 07

The exact values of L_4^r and L_6^r are important to establish what is the role of the strange quark in chiral symmetry breaking, and the m_s -dependence of quantities like F and the chiral condensate – surprises in the m_s expansion? cf. talk by S. Descotes-Genon at Lat07

Nucleon mass

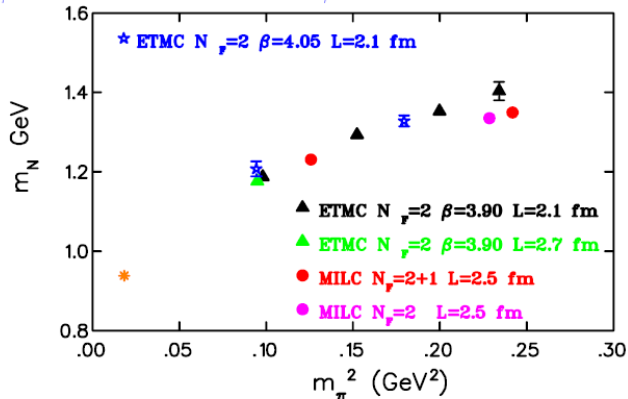
Quark mass dependence known up to $\mathcal{O}(M^4)$:

$$m_N = m - 4c_1 M^2 - \frac{3g^2 M^3}{32\pi F^2} + \left[\tilde{e}_1 - \frac{3(2g^2 - c_2 m)}{8NF^2} \right] M^4 + \mathcal{O}(M^5)$$

The chiral log hidden in \tilde{e}_1 is known, but the LEC e_1^r not

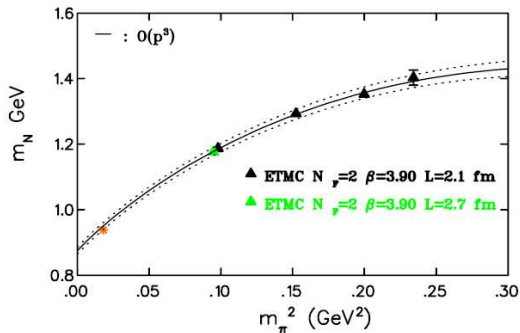
Nucleon mass

Take $a_{\beta=3.9} = 0.0855$ fm and $a_{\beta=4.05} = 0.0666$ fm



Nucleon mass

Take $a = 0.0855$ fm



- Third order:

$$m_N = m_0 - 4c_1 m_\pi^2 - \frac{3g_A^2}{32\pi f_\pi^2} m_\pi^3$$

- Provides a very good fit with $m_0 = 0.875(10)$ GeV and $c_1 = -1.231(17)$ GeV⁻¹ to be compared with $c_1 = -0.9$ GeV⁻¹ extracted from $\pi - N$ -sigma term.
- Physical point, which is not included in the fit, is reproduced.

Chiral expansion of pion form factors

$$\langle \pi^i(p_1) | \bar{q} \Gamma_L q | \pi^j(p_2) \rangle = C_L F_L(s) \quad s = (p_1 + p_2)^2$$

$$L = S \quad \Rightarrow \quad \Gamma_S = 1, \quad C_S = \delta^{ij}$$

$$L = V \quad \Rightarrow \quad \Gamma_V = \tau_3 \gamma_\mu, \quad C_V = i \epsilon^{i3j} (p_1 + p_2)_\mu$$

Taylor expansion near $s = 0$:

$$F_L(s) = F_L(0) \left[1 + \frac{\langle r^2 \rangle_L}{6} s + \mathcal{O}(s^2) \right]$$

Chiral expansion of pion form factors

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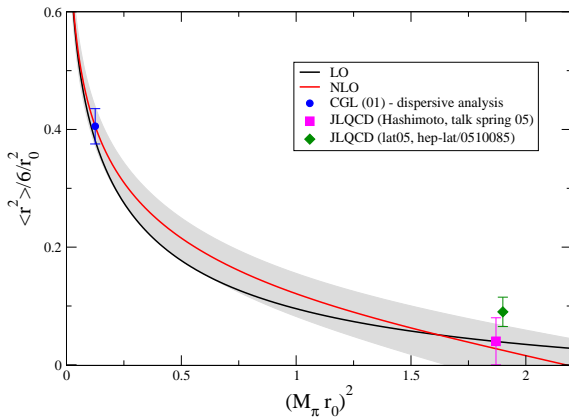
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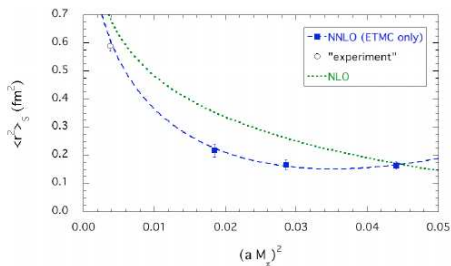
Taylor expansion near $s = 0$:

$$F_L(s) = F_L(0) \left[1 + \frac{\langle r^2 \rangle_L}{6} s + \mathcal{O}(s^2) \right]$$

Chiral expansion of the scalar radius:

$$\langle r^2 \rangle_S = 6 \xi \left\{ \ln \frac{\Lambda_4^2}{M_\pi^2} - \frac{13}{12} - \frac{29}{18} \xi \left(\ln \frac{\Omega_S^2}{M_\pi^2} \right)^2 + \xi c_S + \mathcal{O}(\xi^2) \right\}$$

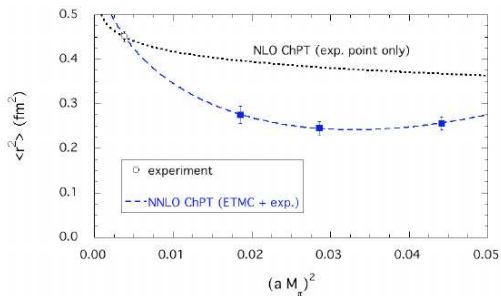
Chiral extrapolation of $\langle r^2 \rangle_{S,V}$ 

Chiral extrapolation of $\langle r^2 \rangle_{S,V}$ 

$$\langle r^2 \rangle_S(NLO) = \frac{12}{(4\pi F)^2} \left[\log\left(\frac{\Lambda_4^2}{M_\pi^2}\right) - \frac{13}{12} \right]$$

fix $aF = 0.0534$ and $\bar{l}_4 = 4.52$ from PLB '07

$$\begin{aligned} \langle r^2 \rangle_S(NNLO) &= \langle r^2 \rangle_S(NLO) \\ &+ d_1 M_\pi^2 + d_2 M_\pi^2 \log(M_\pi^2) \end{aligned}$$

Chiral extrapolation of $\langle r^2 \rangle_{S,V}$ 

$$\langle r^2 \rangle(NLO) = \frac{2}{(4\pi F)^2} \left[\log \left(\frac{\Lambda_6^2}{M_\pi^2} \right) - 1 \right]$$

fix $aF = 0.0534$ from PLB '07

$$\text{exp.} \Rightarrow \bar{\ell}_6 = 14.4 \quad (3)$$

$$\begin{aligned} \langle r^2 \rangle(NNLO) &= \langle r^2 \rangle(NLO) \\ &+ c_1 M_\pi^2 + c_2 M_\pi^2 \log(M_\pi^2) \end{aligned}$$

$$\text{ETMC + exp.} \Rightarrow \bar{\ell}_6 = 17.2 \quad (7)$$

πK vector form factor

In order to extract V_{us} from $K_{\ell 3}$ decays, information on the πK vector form factor at $t = 0$ is needed:

$$f_+^{\pi K}(0) = 1 + f_2 + f_4 + \mathcal{O}(p^6)$$

Ademollo-Gatto theorem $\Rightarrow f_2$ a known function of meson masses:

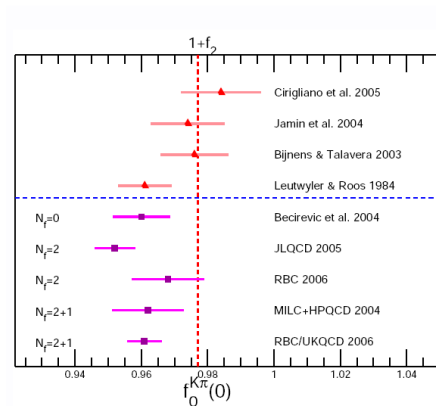
$$f_2 = -0.023$$

Gasser and Leutwyler (85)

$f_+^{\pi K}(0)$ also known to two loops
Problem: determine the relevant LEC's

Post-Schilcher (01), Bijmans-Talavera (03)

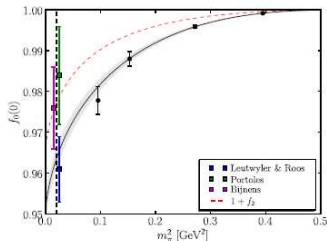
πK vector form factor



- tension between χ PT and lattice?
→ need to reduce error bars
- seems independent of N_f
- in many cases Leutwyler & Roos (1984!) is still used to determine $|V_{us}|$ (cf. PDG)

πK vector form factor

● Better χ^{PT} guidance would help



● Published results cover only 3 of four data points

$$f_+^{K\pi}(0) = 0.9609(51)$$

$$|V_{us}| = 0.2257(9)_{\text{exp}(12)} f_+(0)$$

● Lightest data point will enhance control of chiral limit

● Data set complete ; final analysis & systematic estimates in progress

● Results very much favour Leutwyler-Roos with reduced error

● Big impact on V_{us}

Boyle, talk at Lat07

Accurate chiral extrapolation is essential – two-loop formulae complicated \Rightarrow close collaboration between lattice and chiral people

Nucleons: g_A

Chiral expansion known up to $\mathcal{O}(M^3)$:

$$g_A = g + \left(4\tilde{d}_{16} - \frac{g^3}{NF^2} \right) M^2 + \frac{3(1 + g^2) - 4m(c_3 - 2c_4)}{24\pi mF^2} gM^3 + \mathcal{O}(M^4)$$

Leading chiral logs at $\mathcal{O}(M^4)$ have also been evaluated

Bernard, Meißner (06)

However the chiral expansion behaves badly already for the physical value of the pion mass:

Kambor, Mojžiš (99)

$$\frac{-4m(c_3 - 2c_4)}{24\pi mF^2} M^3 \Big|_{M=0.14\text{GeV}} \sim 12\%$$

Nucleons: g_A

Chiral expansion known up to $\mathcal{O}(M^3)$:

$$g_A = g + \left(4\tilde{d}_{16} - \frac{g^3}{NF^2} \right) M^2 + \frac{3(1+g^2) - 4m(c_3 - 2c_4)}{24\pi mF^2} gM^3 + \mathcal{O}(M^4)$$

Leading chiral logs at $\mathcal{O}(M^4)$ have also been evaluated

Bernard, Meißner (06)

However the chiral expansion behaves badly already for the physical value of the pion mass:

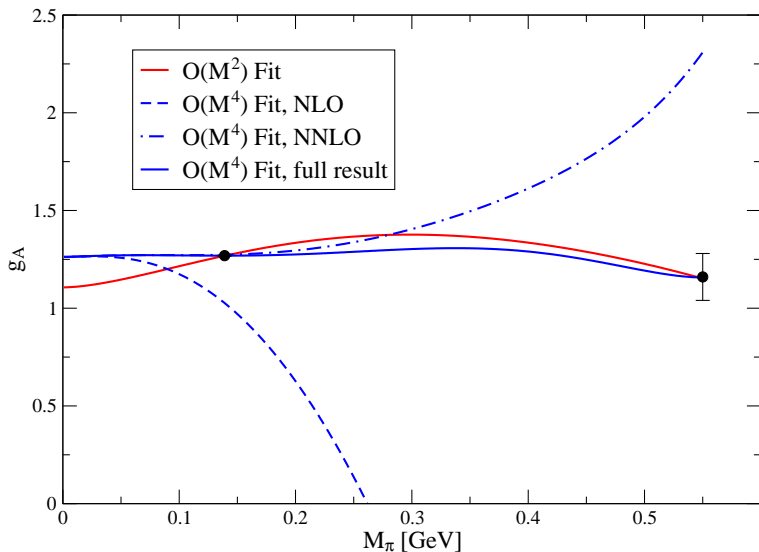
Kambor, Mojžiš (99)

$$\frac{-4m(c_3 - 2c_4)}{24\pi mF^2} M^3 \Big|_{M=0.14\text{GeV}} \sim 12\%$$

Including the Δ as an explicit degree of freedom tames the growth of the M^3 term

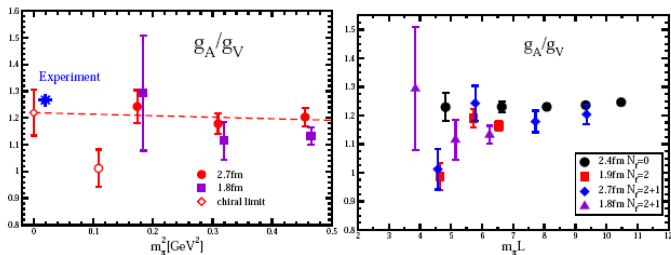
Hemmert, Procura, Weise (03)

But at the price of a fine tuning

Nucleons: g_A 

Nucleons: g_A

- finite volume effect on g_A at lightest mass?
- Chiral limit: $g_A/g_V = 1.220(85)$ omitting lightest mass



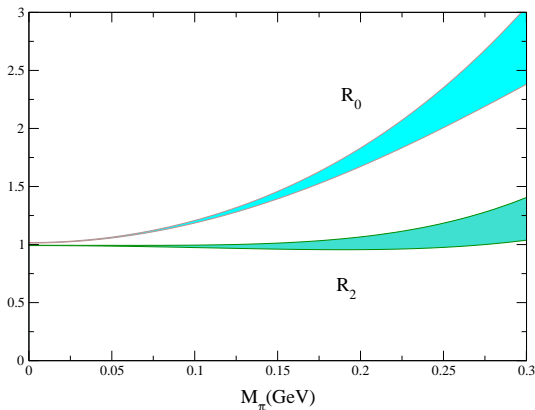
Chiral extrapolation of a_0^0 and a_0^2

Formulae known to two loops

Bijnens, GC, Ecker, Gasser and Sainio (95)

Warning: chiral convergence of the $\pi\pi$ scattering lengths is bad

The relevant mass scale is rather $2M_\pi$!



$$R_l = a_0^l / a_0^l(\text{LO})$$

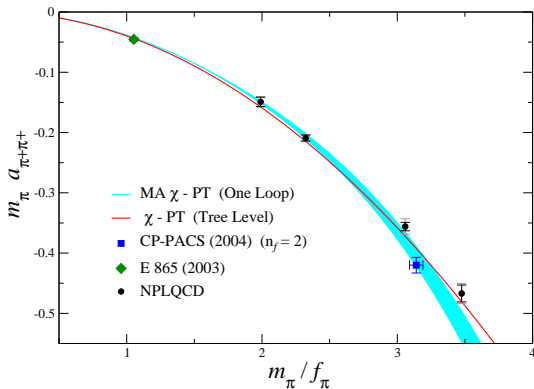
Recent lattice calculations of a_0^2 

Figure from NPLQCD, arXiv:0706.3026

cf. Talk by A. Walker-Loud

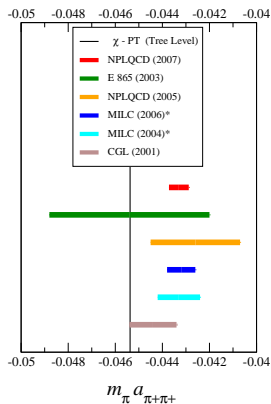
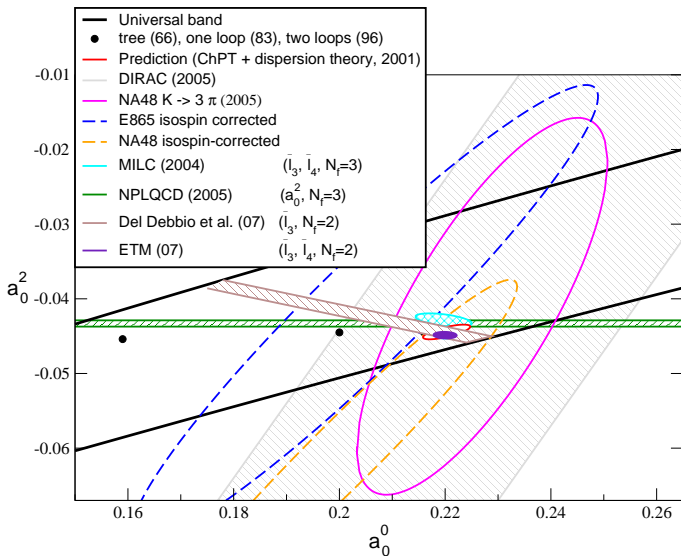
Recent lattice calculations of a_0^2 

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Recent lattice calculations of a_0^2 

Conclusions

- ▶ Recent progress in approaching the chiral region on the lattice is impressive
- ▶ Statistical errors on the low-energy constants determined on the lattice are tiny – **we see already the enormous potential of the method**
- ▶ Study of systematic effects (**continuum extrapolation, finite volume, higher chiral orders**) is only at the beginning
Need to exclude that “the customer is always right”
P.Boyle, Lat07
- ▶ **(Near) future:**
 1. make an effort community-wide to determine as accurately as possible the low-energy constants
 2. use this input in CHPT to describe the phenomenology \Rightarrow test of QCD
- ▶ \Rightarrow ongoing discussion within the Flavianet Lattice Working Group – if you have ideas or proposals, they are most welcome: contact me or Rainer Sommer